1 Second Order Linear ODE: y'' + p(x)y' + q(x)y = r(x)

Solution y = Au(x) + Bv(x) + z(x).

Where u(x), v(x) are Linearly Independent

(i.e. $\forall A, B(\forall x(Au(x) + Bv(x) = 0)) \Rightarrow A = B = 0)$ solutions to the Homogeneous Equation y'' + p(x)y' + q(x)y = 0. Moreover the solution space to the Homogeneous Equation forms a 2 Dimensional Vector Space Over \mathbb{R} and u, v are Basis Vectors which are not unique.

Any solution to the Homogeneous Equation can be written as a Linear Combination of the Basis Vector: i.e. Complementary Solution $y_c = Au(x) + Bv(x)$.

The Particular Solution $y_p = z(x)$ is the constant-less solution to the Non Homogeneous Equation: y'' + p(x)y' + q(x)y = r(x).

The constants A, B can be uniquely determined by the Initial Conditions: y(b), y'(b).

1.1 Wronskian Method

Wronskian
$$W(x) = \left| \begin{pmatrix} u(x) & v(x) \\ u'(x) & v'(x) \end{pmatrix} \right| = u(x)v'(x) - v(x)u'(x)$$

u, v linearly independent is equivalent to $\forall x(W(x) \neq 0)$.

Wronskian satisfies the 1st order homogeneous equation: W'(x) + p(x)W(x) = 0. So we can find W(x) and if we know one solution u(x) to the homogeneous equation we can find the other solution v(x). This result is equivalent to the Reduction of Order method where we assume that v(x) = A(x)u(x) and finding a 1st order ode satisfied by A(x).

Also we can show that the particular solution can be written as z(x) = A(x)u(x) + B(x)v(x). This method is called the Variation of Parameters and we can show that $A(x) = -\frac{v(x)r(x)}{W(x)}$ and $B(x) = \frac{u(x)r(x)}{W(x)}$.

So we can generate the full solution if we know only one solution u(x) to the homogeneous equation.

1.2 Frobenius Method

Consider the 2nd order linear homogeneous ODE, y'' + p(x)y' + q(x)y = 0. Here we try to find at least one solution u(x) as a power series of (x - b). It is convenient to find the constants A, B when the initial conditions y(b), y'(b) are stated at b.

Let $p(x) = \sum_{k=0}^{\infty} p_{k-1}(x-b)^{k-1}$ and $q(x) = \sum_{k=0}^{\infty} p_{k-2}(x-b)^{k-2}$. We are looking for a solution of the form $y = \sum_{k=0}^{\infty} a_k(x-b)^{k+c}$. We need to find both a_k and c. Any coefficients a_k remain indeterminate will become the constants A, B. Here ccan provide powers of (x-b) which are not integer and also help generate both solutions u, v by the same procedure.

We have $y' = \sum_{k=0}^{\infty} a_k (k+c) (x-b)^{k+c-1}$ and $y'' = \sum_{k=0}^{\infty} a_k (k+c) (k+c-1) (x-b)^{k+c-2}$. Substituting into the ODE we have,

$$\sum_{k=0}^{\infty} a_k (k+c) (k+c-1) (x-b)^{k+c-2} + (\sum_{i=0}^{\infty} p_{i-1} (x-b)^{i-1}) (\sum_{j=0}^{\infty} a_j (j+c) (x-b)^{j+c-1})$$

$$+ \left(\sum_{i=0}^{\infty} q_{i-2}(x-b)^{i-2}\right)\left(\sum_{j=0}^{\infty} a_j(x-b)^{j+c}\right) = 0$$

$$\sum_{k=0}^{\infty} a_k (k+c) (k+c-1) (x-b)^{k+c-2} + \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} a_j (j+c) p_{i-1} (x-b)^{i+j+c-2} + \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} a_j q_{i-2} (x-b)^{i+j+c-2} = 0$$

Letting
$$i + j = k$$
 and writing $i = k - j$ we arrive at
 $\sum_{k=0}^{\infty} a_k (k+c)(k+c-1)(x-b)^{k+c-2}$
 $+ \sum_{k=0}^{\infty} \left(\sum_{j=0}^k a_j (j+c) p_{k-j-1} \right) (x-b)^{k+c-2}$
 $+ \sum_{k=0}^{\infty} \left(\sum_{j=0}^k a_j q_{k-j-2} \right) (x-b)^{k+c-2} = 0$
 $\sum_{k=0}^{\infty} \left(a_k (k+c)(k+c-1) + \sum_{j=0}^k a_j ((j+c) p_{k-j-1} + q_{k-j-2}) \right) (x-b)^{k+c-2} = 0$

So we want for all
$$k \ge 0$$

$$0 = a_k(k+c)(k+c-1) + \sum_{j=0}^k a_j((j+c)p_{k-j-1} + q_{k-j-2})$$

$$= a_k((k+c)(k+c-1+p_{-1}) + q_{-2}) + \sum_{j=0}^{k-1} a_j((j+c)p_{k-j-1} + q_{k-j-2})$$

$$= a_k((k+c)^2 + (-1+p_{-1})(k+c) + q_{-2}) + \sum_{j=0}^{k-1} a_j((j+c)p_{k-j-1} + q_{k-j-2})$$

For k = 0 we have $a_0(c^2 + (-1 + p_{-1})c + q_{-2}) = 0$. To keep a_0 as a constant we need c to satisfy the Indicial Equation: $c^2 + (-1 + p_{-1})c + q_{-2} = 0$. Let the roots be α, β . Note that the quadratic equation associated with the coefficient of a_k is $(k + c)^2 + (-1 + p_{-1})(k + c) + q_{-2}$ which has roots $\alpha - k$ and $\beta - k$. The two quadratic equations can have a common root, which occurs when $\beta - k = \alpha$ or when $\alpha - k = \beta$. If $\alpha < \beta$ this occurs when $\alpha + k = \beta$. So this happens when the roots of the indicial equation is differ by an integer and when we put $c = \alpha$ which is the smaller root. We can think about 3 cases

Case 1: $\alpha \neq \beta$ and are not differ by an integer In this case a_k can be expressed by a_0 only. We can have two independent solutions by substituting $c = \alpha, \beta$.

Case 2: $\alpha \neq \beta$ and are differ by an integer k = n.

In this case a_n can be expressed by a_0 only for the larger root $c = \beta$ and we get one solution.

But we cannot express a_n by a_0 for the smaller root $c = \alpha$. So if the equations does not force $a_0 = 0$, then all coefficients are expressible by a_0 and a_n . Which means you end up with two solutions followed by two constants a_0, a_n if the second one is not merely a shift of index of the first one.

But if the equations force $a_0 = 0$ you end up expressing all coefficients by a_n which means you get one solution which can still be a shift of index to the $c = \alpha$ case.

Case 3: $\alpha = \beta$

In this case coefficient of a_k will never be 0. So close to $c = \alpha$ we can force $a_k((k+c)^2 + (-1+p_{-1})(k+c) + q_{-2}) + \sum_{j=0}^{k-1} a_j((j+c)p_{k-j-1} + q_{k-j-2}) = 0$ for all $k \ge 1$ and all coefficients can be expressible by a_0 and c. We end up with a different differential equation

 $y'' + p(x)y' + q(x)y = a_0(c - \alpha)^2 x^{c-2}$ with y = y(x, c) be a solution. Note that $y = y(x, \alpha)$ is a solution. We will show that $y_c(x, \alpha)$ is the second independent solution.

In this context $\frac{d}{dx} = \frac{\partial}{\partial x}$ and differentiating both sides partially wrt c and with $y \in \mathcal{C}^3$ we have: $\frac{\partial}{\partial c} \frac{\partial^2 y}{\partial x^2} + p(x) \frac{\partial}{\partial c} \frac{\partial y}{\partial x} + q(x) \frac{\partial y}{\partial c} = \frac{\partial^2}{\partial x^2} \left(\frac{\partial y}{\partial c}\right) + p(x) \frac{\partial}{\partial x} \left(\frac{\partial y}{\partial c}\right) + q(x) \left(\frac{\partial y}{\partial c}\right) = 2a_0(c-\alpha)x^{c-2} + a_0(c-\alpha)^2x^{c-2}\log x$ Now putting $c = \alpha$ we have: $\frac{\partial^2}{\partial x^2} \left(\frac{\partial y}{\partial c}|_{c=\alpha}\right) + p(x)\frac{\partial}{\partial x} \left(\frac{\partial y}{\partial c}|_{c=\alpha}\right) + q(x) \left(\frac{\partial y}{\partial c}|_{c=\alpha}\right) = 0.$ This means $\frac{\partial y}{\partial c}|_{c=\alpha} = y_c(x,\alpha)$ is the second solution.

1.3 Radius of Convergence

Power series are functions of the form $f(x) = \sum_{k=0}^{\infty} a_k (x-b)^k$. They converge Absolutely and Uniformly for |x-b| < R and diverge for |x-b| > R for some $R \ge 0$. R is called the radius of convergence and we can use the Root test to find it: $R = \frac{1}{\lim_{k \to \infty} |a_k|^{1/k}}$. Let f, g are power series with radii of convergence R_f, R_g . f + g is a power series $(f + g)(x) = \sum_{k=0}^{\infty} A_k (x-b)^k$ where $A_k = a_k + b_k$. fg is a power series defined as the Cauchy Product $(fg)(x) = \sum_{k=0}^{\infty} A_k (x-b)^k$ where $A_k = \sum_{j=0}^k a_j b_{k-j}$. Both the sum and product power series converges with a radius of convergence min (R_f, R_g) . With $p(x) = \sum_{k=0}^{\infty} p_{k-1}(x-b)^{k-1}$ and $q(x) = \sum_{k=0}^{\infty} p_{k-2}(x-b)^{k-2}$ note that (x-b)p(x) and $(x-b)^2q(x)$ have power series and have radius of convergence R_p and R_q respectively. Note that radii of convergence of y, y', y'' are all equal and be R. Since we are looking for solutions to y'' + p(x)y' + q(x)y = 0 or y'' = -p(x)y' - q(x)y

we have: $R = \min(\min(R_p, R), \min(R_q, R))$. $R = \min(R_p, R_q)$ satisfies this equation.

1.4 Legendre Equation

Legendre ODE $(1 - x^2)y'' - 2xy' + n(n+1)y = 0$ Solution: $y = AP_n(x) + BQ_n(x)$

 $P_n(x)$:Legendre functions of the first kind(Legendre polynomials when $n \in \mathbb{Z}^+$) $Q_n(x)$:Legendre functions of the second kind

 $P_n(x)$ satisfies the Rodrigue's formula: $P_n(x) = \frac{1}{2^n n!} \frac{d^n}{dx^n} \left((x^2 - 1)^n \right)$

For $f,g \in \mathcal{C}[-1,1]$ with the Inner Product $\langle f,g \rangle = \int_{-1}^{1} f(x)g(x)dx$, Legendre polynomials satisfies the Orthogonality Condition: $\int_{-1}^{1} P_m(x)P_n(x)dx = 0$ if $m \neq n$ and $\frac{2}{2n+1}$ if m = n.

This allows us to write the Legendre Series: $f(x) = \sum_{k=0}^{\infty} a_k P_k(x)$ where $a_k = \frac{2k+1}{2} \int_{-1}^{1} f(x) P_k(x) dx$.

1.5 Associated Legendre Equation

Associated Legendre ODE $(1 - x^2)y'' - 2xy' + \left(n(n+1) - \frac{m^2}{1-x^2}\right)y = 0$ Solution: $y = AP_n^m(x) + BQ_n^m(x)$ $P_n^m(x) = (1-x^2)^{m/2}\frac{d^m}{dx^m}P_n(x)$:Associated Legendre functions of the first kind(Associated Legendre polynomials when $n \in \mathbb{Z}^+$) $Q_n(x) = (1 - x^2)^{m/2}\frac{d^m}{dx^m}Q_n(x)$:Associated Legendre functions of the second kind. Also see the Mathematica file: Legendre.nb

1.6 Gamma Function and Pochhammer Symbol

Gamma Function: $\Gamma(x) = \int_0^\infty e^{-t} t^{x-1} dt, x > 0.$ Gamma function satisfies use the formula $\Gamma(x+1) = x\Gamma(x)$ which implies $n! = \Gamma(n+1)$ for a none negative integer n.

We define the Gamma function for x < 0 by $\Gamma(x) = \frac{\Gamma(x+1)}{x}$.

Since $\Gamma(x) \to \infty$ as $x \to 0^+$ we have $|\Gamma(x)| \to \infty$ as $x \to -n$ where $n \in \mathbb{Z}^+$. Pochhammer symbol: $(\alpha)_k = \alpha(\alpha + 1) \cdots (\alpha + k - 1)$. Note that $(1)_n = n!$ for none negative integer n. We can show that $(\alpha)_k = \frac{\Gamma(\alpha+k)}{\Gamma(\alpha)}$.

1.7 Bessel Equation

Bessel ODE $x^2y'' + xy' + (x^2 - \nu^2)y = 0$ Solution: $y = AJ_{\nu}(x) + BY_{\nu}(x)$ $J_{\nu}(x) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k!\Gamma(k+1+\nu)} \left(\frac{x}{2}\right)^{2k+\nu}$: Bessel Function of the First Kind If $\nu \notin \mathbb{Z}$ the solution is $y = AJ_{\nu}(x) + BJ_{-\nu}(x)$. When $\nu = n \in \mathbb{Z}$ we have $J_n(x) = (-1)^n J_{-n}(x)$, so $J_{-n}(x)$ is not the second independent solution. Note that $\cos n\pi = (-1)^n$ for $n \in \mathbb{Z}$. Motivated by this we define a second solution for $\nu \notin \mathbb{Z}$: $Y_{\nu}(x) = \frac{J_{\nu}(x)\cos(\pi\nu) - J_{-\nu}(x)}{\sin(\pi\nu)}, \nu \in \mathbb{Z}$: Bessel Function of the second Kind When $\nu = n \in \mathbb{Z}$ we use the L'Hopital Rule and define: $Y_n(x) = \frac{1}{\pi} \left(\frac{\partial}{\partial\nu} J_{\nu}(x)|_{\nu=n} - (-1)^n \frac{\partial}{\partial\nu} J_{-\nu}(x)|_{\nu=n}\right)$. Let $x = \lambda_k, k \ge 0$ be roots of $J_{\nu}(ax) = 0$. For $f, g \in \mathcal{C}^1[0, a]$ with the Inner Product $< f, g \ge \int_0^a xf(x)g(x)dx$, Bessel Functions satisfies the Orthogonality Condition: $\int_0^a xJ_{\nu}(\lambda_m x)J_{\nu}(\lambda_n x)dx = 0$ if $m \ne n$ and $\frac{a^2}{2} \left(J'_{\nu}(\lambda_n a)\right)^2$ when m = n. This allows us to write the Bessel Series: $f(x) = \sum_{k=0}^{\infty} a_k J_{\nu}(\lambda_k x)$ where $a_k = \frac{2}{a^2(J'_{\nu}(\lambda_k a))^2} \int_0^a xf(x)J_{\nu}(\lambda_k x)dx$. Also see the Mathematica file: Bessel.nb

Definition 1. L^p spaces and norms

 $||f||_{p} = \left(\int_{-\infty}^{\infty} |f(t)|^{p} dt\right)^{\frac{1}{p}}, 1 \le p < \infty, \ L^{p}(\mathbb{R}) = \{f : ||f||_{p} < \infty\}$ $||f||_{\infty} = \sup\{|f(x)| : x \in \mathbb{R}\}, \ L^{\infty}(\mathbb{R}) = \{f : ||f||_{\infty} < \infty\}$

Definition 2. Schwartz spaces and norms $\|f\|_{m,n} = \sup\{|x^m f^n(x)| : x \in \mathbb{R}\}\$ $\mathcal{S} = \{f \in \mathcal{C}^\infty : \forall m, n \in \mathbb{N}, \|f\|_{m,n} < \infty\}$ Note that if $f \in \mathcal{S}$ then $x^m f^n(x) \to 0$ as $x \to \pm \infty$ for all $m, n \in \mathbb{N}$

Theorem 1. $L^2 \subset L^1 \subset L^\infty$ $S \subset L^p, 1 \le p \le \infty$

Definition 3.

Fourier Transform of $f: \mathcal{F}(f)(\omega) = \int_{-\infty}^{\infty} f(t)e^{-i\omega t}dt$ Inverse Fourier Transform of $F: \mathcal{F}^{-1}(F)(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega)e^{-i\omega t}d\omega$

Theorem 2.

If $f \in L^1(\mathbb{R}) \Rightarrow \mathcal{F}(f)(\omega)$ exists, continuous and $\to 0$ as $\omega \to \pm \infty$ If $f \in S$ then $\mathcal{D}(f) = f' \in S$ If $f \in S$ then $\mathcal{F}(f) \in S$ and $\mathcal{F}^{-1}\mathcal{F}(f) = f$ so $\mathcal{F} : S \to S$ is a bijection If $F \in S$ then $\mathcal{F}^{-1}(F) \in S$ and $\mathcal{F}\mathcal{F}^{-1}(F) = F$ so $\mathcal{F}^{-1} : S \to S$ is a bijection

Definition 4.

Inner Product, $\langle f, g \rangle = \int_{-\infty}^{\infty} f(x)\overline{g(x)}dx$. So $||f||_2 = \sqrt{\langle f, f \rangle}$ $(f,g) = \int_{-\infty}^{\infty} f(x)g(x)dx$. So $(f,g) = \langle f, \overline{g} \rangle$ Convolution, $(f * g)(t) = \int_{-\infty}^{\infty} f(x)g(t-x)dx$

Theorem 3. $f, g \in S$

$$\begin{split} \int_{-\infty}^{\infty} f'(x)g(x)dx &= -\int_{-\infty}^{\infty} f(x)g'(x)dx, or (f',g) = -(f,g') \\ (\mathcal{I}f,g) &= (\mathcal{F}f)(0)(\mathcal{F}g)(0) - (f,\mathcal{I}g), \ (\mathcal{I}f)(x) = \int_{-\infty}^{x} f(t)dt \\ \mathcal{F}(e^{iat}f(t))(\omega) &= i\omega\mathcal{F}(f)(\omega - a) \\ \mathcal{F}(f')(\omega) &= i\omega\mathcal{F}(f)(\omega) \\ \mathcal{F}(f'')(\omega) &= -\omega^{2}\mathcal{F}(f)(\omega) \\ \int_{-\infty}^{\infty} \mathcal{F}f(x)g(x)dx = \int_{-\infty}^{\infty} f(x)\mathcal{F}g(x)dx, \ or \ (\mathcal{F}f,g) = (f,\mathcal{F}g) \\ \int_{-\infty}^{\infty} \mathcal{F}^{-1}f(x)g(x)dx = \int_{-\infty}^{\infty} f(x)\mathcal{F}^{-1}g(x)dx, \ or \ (\mathcal{F}^{-1}f,g) = (f,\mathcal{F}^{-1}g) \\ \int_{-\infty}^{\infty} |f(x)|^{2} &= \frac{1}{2\pi}\int_{-\infty}^{\infty} |\mathcal{F}f(x)|^{2}dx \ or \ ||f||_{2} &= \frac{1}{\sqrt{2\pi}}||\mathcal{F}f||_{2} \end{split}$$

Note 1. We are going to take $f : \mathbb{R} \to \mathbb{C}$ in S and going to convert it into a function $f : S \to \mathbb{C}$ by $f(g) = (f,g) = \int_{-\infty}^{\infty} f(x)g(x)dx$ for $g \in S$. We note that 1.f is continuous. Equivalently if $g_n \to g$ in S then $f(g_n) \to f(g)$ i.e. $(f,g_n) \to (f,g)$ 2.f is linear. Equivalently if $g, h \in S$ and $a, b \in \mathbb{C}$ then f(ag + bh) = af(g) + bf(h) i.e. (f, ag + bh) = a(f,g) + b(f,h)

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Now consider the four operators defined for $f \in S$: Differential operator $\mathcal{D} : (\mathcal{D}f)(x) = f'(x)$ Fourier Transform \mathcal{F} : Inverse Fourier Transform \mathcal{F}^{-1} : Integral operator $\mathcal{I} : (\mathcal{I}f)(x) = \int_{-\infty}^{x} f(t)dt$ It can be shown that all the above operations except the last one land on S. In other words if $f \in S$ then $\mathcal{D}f, \mathcal{F}f, \mathcal{F}^{-1}f \in S$. We also have noted that all the following are true 3. $(\mathcal{D}f)(g) = -f(\mathcal{D}g)$ i.e. $(\mathcal{D}f,g) = -(f,\mathcal{D}g)$ 4. $(\mathcal{F}f)(g) = f(\mathcal{F}g)$ i.e. $(\mathcal{F}f,g) = (f,\mathcal{F}g)$ 5. $(\mathcal{F}^{-1}f)(g) = f(\mathcal{F}^{-1}g)$ i.e. $(\mathcal{F}^{-1}f,g) = (f,\mathcal{F}^{-1}g)$ 6. $(\mathcal{I}f)(g) = (\mathcal{F}f)(0)(\mathcal{F}g)(0) - f(\mathcal{I}g)$ i.e. $(\mathcal{I}f,g) = (\mathcal{F}f)(0)(\mathcal{F}g)(0) - (f,\mathcal{I}g)$ Now we are going to extend the class of functions S to include "functions" which are not cooked out of $f \in S$ as in the above process. Definition S a Termeered Distribution $f \in \mathcal{T}$ is a function $f : S \to C$

Definition 5. a Tempered Distribution $f \in \mathcal{T}$ is a function $f : S \to \mathbb{C}$.

Satisfying all the above properties 1 and 2: Continuous and Linear.

We see form the above discussion that $\mathcal{S} \subset \mathcal{T}$.

For $g \in S$, we will still write f(g) = (f, g) as necessary although there is no (f, g) realized as an integral

We can use the formulas 3 through 6 to define the operations $\mathcal{D}, \mathcal{F}, \mathcal{F}^{-1}, \mathcal{I}$ on \mathcal{T} . It can be also shown now that all the 4 operations land on \mathcal{T} . In other words if $f \in \mathcal{T}$ then $\mathcal{D}f, \mathcal{F}f, \mathcal{F}^{-1}f, \mathcal{I}f \in \mathcal{T}$ so $\mathcal{D}, \mathcal{F}, \mathcal{F}^{-1}, \mathcal{I}: \mathcal{T} \to \mathcal{T}$.

Definition 6.

Delta function $\delta \in \mathcal{T}$ and $\delta(f) = (\delta, f) = f(0)$ Heaviside function H(t) = 1 if $t \ge 0$ and 0 otherwise Rectangle Function $\Pi(t) = 1$ if $-\frac{1}{2} \le t \le \frac{1}{2}$ and 0 otherwise. Sinc Function sincx $= \frac{\sin x}{x}$

Theorem 4.

 $\mathcal{F}\delta = 1$ $\mathcal{F}1 = 2\pi\delta$ $\mathcal{D}H = \delta$ $\mathcal{I}\delta = H$ $\mathcal{F}(e^{iat})(\omega) = 2\pi\delta(\omega - a)$ $\mathcal{F}(\sin at)(\omega) = \frac{\pi}{i} (\delta(\omega - a) - \delta(\omega + a))$ $\mathcal{F}(\cos at)(\omega) = \pi (\delta(\omega - a) + \delta(\omega + a))$ $\mathcal{F}(\Pi)(\omega) = sinc(\omega/2)$ $\mathcal{F}(f * g)(\omega) = \mathcal{F}(f)(\omega)\mathcal{F}(g)(\omega)$

Example 1. Use Fourier Transform to Solve the Laplace's equation $\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$. u(x,0) = f(x), u(x,y) bounded as $y \to \infty$. Compute the answer for f(x) = x for 0 < x < 1 and 0 otherwise.

Do the same question with the given boundary condition replaced by $\frac{\partial u}{\partial x}(x,0) = f(x)$